

This is the preprint of the contribution published as:

Shan, Y., Liu, L., Liu, Y., Harms, H., Wick, L.Y. (2020):

Effects of electrokinetic phenomena on bacterial deposition monitored by quartz crystal microbalance with dissipation monitoring

Environ. Sci. Technol. **54** (21), 14036 – 14045

The publisher's version is available at:

<http://dx.doi.org/10.1021/acs.est.0c04347>

1 **Effects of Electrokinetic Phenomena on Bacterial Deposition monitored by Quartz**
2 **Crystal Microbalance with Dissipation Monitoring (QCM-D)**

3

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16 Intended for: Environmental Science & Technology

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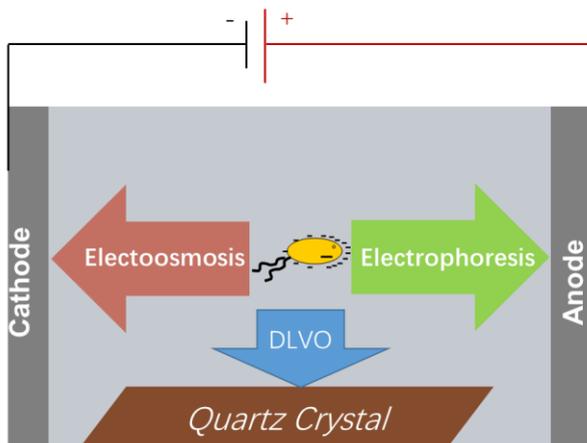
27 **Abstract**

28 Bacterial deposition is the first step in the formation of microbial biofilms in environmental
29 technology, and there is high interest in controlling such deposition. Earlier work indicated that
30 direct electric current (DC) fields could influence bacterial deposition in percolation columns.
31 Here, a time-resolved quartz crystal microbalance with dissipation monitoring (QCM-D) and
32 microscopy-based cell counting were used to quantify DC field effects on the deposition of
33 bacterial strains *Pseudomonas putida* KT2440 and *Pseudomonas fluorescens* LP6a at varying
34 electrolyte concentrations and weak electric field strengths (0-2 V cm⁻¹). DC-induced
35 frequency (Δf) shifts, dissipation energy (ΔD), and ratios thereof ($\Delta f/\Delta D$) proved as good
36 indicators of the rigidity of cell attachment. We interpreted QCM-D signals using a theoretical
37 approach calculating the attractive DLVO-force and the shear and drag forces acting on a
38 bacterium near collector surfaces in a DC electric field. We found that changes in DC-induced
39 deposition of bacteria depended on the relative strengths of electrophoretic drag and
40 electroosmotic shear forces. This could enable the prediction and electrokinetic control of
41 microbial deposition on surfaces in natural and manmade ecosystems.

42

43 **Keywords:** bacterial deposition, DLVO, DC electric fields, electrokinetics, electroosmosis,
44 electrophoresis.

45 **Abstract Art**



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48

49 **Introduction**

50 Microbial biofilms provide essential ecosystem services in many natural and manmade
51 environments. While being beneficial in e.g. wastewater treatment systems or the degradation
52 of contaminants, biofilms can also be detrimental to both human health and industrial
53 applications. Biofouling can increase the corrosion of metals,¹ infect medical devices,^{2,3} and
54 pollute drinking water systems.⁴⁻⁶ Direct current (DC) electric fields and their associated
55 electrokinetic phenomena have been found to affect the bacterial deposition⁷⁻¹¹ that precedes
56 biofilm formation. DC electric fields evoke various electrokinetic transport processes in both
57 conductive^{12,13} and non-conductive matrices^{14,15} immersed in liquid. Electric field applied in
58 the liquid surrounding non-conductive materials may induce electrokinetic phenomena, which
59 allow for targeted movement of bacteria and colloidal particles in the system, even in the
60 absence of pressure-driven hydraulic flow.^{10,19-21} While electromigration and electrophoresis
61 refer to the transport of charged molecules and particles to the electrode of opposite charge,
62 electroosmosis reflects the surface charge-induced movement of pore fluids, usually from the
63 anode to the cathode (electroosmotic flow, EOF).²² Due to a plug-shaped flow profile that acts
64 a few nanometres above a surface, EOF is thought to affect bacterial deposition by inducing
65 shear forces (F_{EOF}).²³⁻²⁵ Electrophoresis (EP), by contrast, induces a drag force (F_{EP}) on the
66 (negatively) charged bacteria²⁶⁻²⁸ and hence acts in the direction opposite to F_{EOF} . A bacterium
67 approaching a surface or being located at a distance of the secondary DLVO energy minimum²⁹)
68 will be subject to F_{EOF} and F_{EP} and the relative strength of the two forces has been proposed to
69 be a driver for observed DC field effects on bacterial deposition.^{14,26,30,31} Electrokinetic
70 phenomena are directly correlated to the electric field strength (E) applied, the surface
71 properties of the matrices and the (bio-)colloidal particles, and the ionic strength of the
72 electrolytes; i.e. parameters that may impact interactions between bacterium and solid
73 surfaces.^{32,33} Here we assessed the effect of DC electric fields on bacterial deposition using a

74 quartz crystal microbalance with dissipation (QCM-D) that allows for real-time
75 characterization of bacteria-surface interactions^{34,35} and, hence, also electrokinetic effects on
76 bacterial deposition during transport in porous media. QCM-D reflects the amount and
77 viscoelastic properties of an adhering mass (bacteria) by changes in the resonance frequency
78 (Δf) and changes in the energy dissipation (ΔD) of an oscillating crystal coating sensor
79 surface.³⁶⁻³⁹ The Δf is an indicator of the bacterial mass attached to the sensor while ΔD
80 indicates the softness of non-rigid mass adhesion.^{40,41} Given constant temperature, liquid
81 viscosity and density, and flow velocity both signals vary according to the surface charge and
82 the hydrophobic properties of the bacteria and the sensor surface during the monitoring of
83 bacteria-surface interactions.⁴²⁻⁴⁴ A plot of ΔD versus Δf compares the induced energy
84 dissipation per coupled unit mass: lower $\Delta f/\Delta D$ values indicate the formation of a dissipative,
85 soft, and fluid film, while higher $\Delta f/\Delta D$ values suggest a more rigid layer of attached bacterial
86 mass.^{34,45} Hence, the Δf and ΔD of the QCM-D sensor allow to analyze the diverse responses
87 and transition from inertial to elastic loading of cells having similar surface morphologies in
88 the presence and absence of external electric fields, and hence allow to deduce the mechanisms
89 of electrokinetic effects on the surface-bacteria bond.^{46,47} If Δf values are supported by direct
90 microscopy observed cell density, QCM-D monitoring can be used to quantify the rate of
91 bacterial attachment to the sensor surface, to approximate the time-resolved electrokinetic
92 effects on bacterial deposition at varying environmental conditions, and to compare bacterial
93 deposition to electrokinetically induced forces (F_{EOF} and F_{EP}) acting on bacteria adjacent to a
94 solid collector surface.

95 Here we used a QCM-D approach to assess the joint effects of a DC electric field and the ionic
96 strength of the electrolyte on the deposition at a nanogram level of two bacteria of differing
97 physicochemical cell surface properties and opposite transport behaviour in percolation
98 columns exposed to external DC fields.^{14,48} QCM-D data were supported by microscopic cell

99 counting and analyzed by a recently published theoretical approach that involved calculating
100 the DLVO colloidal interaction, the hydraulic drag, and the electrokinetic forces acting on a
101 bacterium near a collector surface in a DC electric field.

102

103 **Materials and Methods**

104 **Cultivation of bacteria and inoculum preparation**

105 *Pseudomonas putida* KT2440 (GenBank accession No. AE015451)⁴⁹ and *Pseudomonas*
106 *fluorescens* LP6a (GenBank accession No. AF525494)⁵⁰ were cultivated in minimal media
107 with 1.0 gL⁻¹ glucose as a carbon source until the early stationary phase (25 °C; rotary shaker
108 at 150 rpm). The cultures were then centrifuged at 3000 × g and resuspended in 10 mM (5
109 mmol K₂HPO₄ and 5 mmol KH₂PO₄ diluted in 1 L deionized water), 50 mM (29 mmol K₂HPO₄
110 and 21 mmol KH₂PO₄ diluted in 1 L DI water), and 100 mM (61 mmol K₂HPO₄ and 39 mmol
111 KH₂PO₄ diluted in 1 L DI water) potassium phosphate buffer, pH = 7 (PB) using a Vortex
112 mixer (Vortex-Genie 2, Scientific Industries, USA) to obtain bacterial suspensions with an
113 optical density at 600 nm of 0.30 (OD_{600 nm} = 0.30).

114

115 **Characterization of physiochemical properties of bacterial and sensor surfaces**

116 The zeta-potentials of bacteria (ζ_{bac}) and silica beads (ζ_{s}) were measured by Doppler
117 electrophoretic light scattering analysis (Zetasizer Nano ZS, Malvern Instruments, Malvern,
118 UK) with a Dip Cell Kit. The zeta potential of the silica sensor surface was estimated using
119 smashed silica beads in the different electrolytes. Clean glass beads were smashed with a
120 mortar and a pestle to a size of < 100 μm, heated at 200 °C in a muffle furnace for 2 h, then
121 allowed to cool to room temperature (25 °C) under sterile conditions. The contact angles (θ) of
122 the bacterial strains and the sensor were quantified using a DSA 100 drop-shape analysis
123 system (Krüss GmbH, Hamburg, Germany) in three solvents (water, formamide, methylene

124 iodide)^{15,51} and are listed in Table S1. Bacterial lawns were prepared by depositing bacteria
125 from inoculated suspensions on cellulose acetate membrane filters (Millipore, 0.45 μm); four
126 droplets were applied per filter, in triplicate experiments for each solvent.

127

128 **QCM-D analysis of cell deposition on the silica sensor surface**

129 Interactions between bacterial cells and a silica surface were studied with an E4 QCM-D unit
130 (Q-Sense AB, Gothenburg, Sweden) using silica-coated sensor chips (QSX-303, 5 MHz, AT-
131 cut, diameter: 14 mm, Q-Sense AB, Gothenburg, Sweden). Experiments were performed in a
132 QCM-D system comprised of an inlet solution container, four QCM-D chambers, a buffering
133 bottle, and a wastewater container (for a schematic view of the set-up cf. Fig. S1). Bacterial
134 suspensions were pumped through QCM-D tubing under pressure-driven flow using a digital
135 peristaltic pump (ISM932A, Ismatec, Cole-Parmer, Canada) at a fixed flow rate of 200 $\mu\text{L min}^{-1}$
136 (flow velocity: $6 \times 10^{-7} \text{ m s}^{-1}$) at $20 \pm 0.2 \text{ }^\circ\text{C}$ (cf. Fig. S1). DC fields ($E = 0.5, 1.0, \text{ and } 2.0 \text{ V}$
137 cm^{-1}) were generated by a power pack (BK Precision 9174), and connected to two Ti/Ir
138 electrodes placed in the bacterial suspension (cathode) and the anode bottle. As extensions of
139 the electrodes, two copper wires (0.2 mm i.d., renewed after each experiment) were connected
140 to the Ti/Ir electrodes. They were cautiously inserted into the tubing up to a distance of 2 mm
141 from the inlet and outlet of the QCM-D chamber, resp.; i.e. with no contact to the sensor.
142 Placement of the anode wire outside the QCM-D chamber avoided possible interferences of
143 electrochemically released copper ions with the QCM-D measurements. Placing the copper
144 electrodes wires close to the QCM-D chamber allowed us to apply low potential while
145 simultaneously maintaining the electric field strength in the QCM-D chamber as detailed below.
146 PB at either 10, 50, or 100 mM was used as the electrolyte and DC electric fields of $E = 0, 0.5,$
147 $1.0, \text{ or } 2.0 \text{ V cm}^{-1}$ were applied. Prior to the experiment, clean sterilized silica sensors were
148 mounted in the QCM-D chamber, and the screws on the back of QCM-D chambers were sealed

149 until hand-tight, then locked by the snap on the base bracket. The frequency and dissipation of
150 silica sensors in DI water were assured to deviate less than $\pm 10\%$ from the standard frequency
151 and dissipation values at overtones 1, 3, 5, 7, 9, 11, and 13 (corresponding to 5, 15, 25, 35, 45,
152 55, and 65 MHz), respectively. Identical ΔD and Δf signals were detected in controls pumping
153 cell-free buffer solutions in presence and absence of external DC fields. Before proceeding
154 with experiments, the system baseline was stabilized by pumping through ultrapure water for
155 20 min, followed by cell-free PB (of ionic strength equal to that of the cell suspensions) for 40
156 min. Bacterial suspensions of either *P. putida* KT2440 or *P. fluorescens* LP6a (in 10, 50, or
157 100 mM PB) were then pumped into the QCM-D chamber over 2 hours and the frequency and
158 dissipation was monitored simultaneously. Experiments were performed in triplicate at $E = 0$,
159 0.5, 1.0, and 2.0 V cm⁻¹.

160 After each experiment, the sensors were rinsed with 1.5 mL ultrapure water in a 50 mL
161 centrifuge tube and bacterial cells were detached using an ultrasonic washing unit (FS60, Fisher
162 Scientific, Canada) for 10 min. The sensor was removed using tweezers, disinfected in a UV
163 chamber for 20 min, cleaned in 50 mL of 2% sodium dodecyl sulfate (SDS), rinsed thoroughly
164 with ultrapure water, dried under a nitrogen stream, and sterilized for 20 min in a UV chamber
165 following the washing protocol provided with the silica sensors.

166

167 **Microscopic quantification of cells attached to the sensor**

168 At the end of each QCM-D analysis (i.e., after 2 h) the bacteria the cells on the sensor were
169 detached with an ultrasonic unit for 10 minutes and collected in 1.5 mL water. Detachment was
170 complete as verified by microscopic analysis of the sensor surface. The bacterial suspension
171 was centrifuged at $6000 \times g$ for 5 min, then 1.45 mL of the supernatant was removed. The
172 bacterial pellet was resuspended in the residual liquid (0.05 mL) with a Vortex mixer (Vortex-
173 Genie 2, Cole-Parmer, Canada). The suspension was then injected to a Hemacytometer

174 (Improved Neubauer 0.1 mm, Fisher Scientific, Canada) to take pictures and quantify the
175 bacterial cell concentration by epifluorescence microscopy (Axioskop II microscope, Carl
176 Zeiss, Canada) equipped with a camera (Carl Zeiss Microimaging GmbH, Canada). Pre-
177 experiments were conducted to observe the distribution of attached cells on the whole sensor
178 surface and the efficiency of ultrasonic cell detachment, resp. Images were analyzed by ImageJ
179 software (ImageJ 1.46r, USA) to quantify the cells. The automatic counting codes used for cell
180 counting are listed in the supporting information. The density of the cells removed from the
181 sensor surface (d_c) was calculated by dividing the number of cells detached from each sensor
182 by the sensor surface area.

183

184

185 **Theory**

186 **Forces acting on bacteria on a collector surface**

187 Although the Derjaguin, Landau, Verwey, and Overbeek (DLVO) theory of colloidal
188 interactions⁵²⁻⁵⁴ does not account for surface heterogeneity, hydration effects, or hydrophobic
189 interactions, it is a powerful predictor of bacterial deposition in solutions of high ionic strength
190 ($I = 0.1-0.3$ M).^{29,55-57} DLVO interaction energy profiles of bacterial deposition depend on the
191 physicochemical properties of the microbe, the collector surface, and the ionic strength of the
192 aqueous medium. DLVO theory also predicts reversible bacterial deposition even at high
193 attractive forces^{58,59} at a so-called secondary minimum of the energy profile, typically located
194 5-20 nm above a collector surface. Therefore, net forces acting on bacteria in the secondary
195 minimum may influence bacterial deposition, bacterial attachment, and biofilm formation.^{14,15}
196 The net force at the secondary minimum is estimated to act on a bacterium through a
197 combination of the DLVO force of colloidal interaction (F_{DLVO}), the hydraulic flow shear force

198 (F_{HF}), the electroosmotic flow shear force (F_{EOF}), and the electrophoretic drag force (F_{EP}),¹⁴ as
 199 shown in Eq. 1:

$$200 \quad F_{net} = F_{DLVO} + F_{HF} + F_{EOF} + F_{EP}. \quad (1)$$

201 The DLVO interaction force and hydraulic force are depicted in Eqs. S1-S11. It should be noted
 202 that the DLVO force is calculated at the secondary minimum distance, where the DLVO
 203 interaction controls the reversible bacterial deposition.⁵⁹ The electroosmotic shear force can be
 204 calculated with Eq. 2:

$$205 \quad F_{EOF} = F_d^* \times 6\pi\eta a V_{EOF} = F_d^* \times 6\pi\eta a \times \left[-\frac{\varepsilon_0 \varepsilon_r \zeta_s E}{\eta} \left(1 - \frac{2I_1(\kappa h_s)}{\kappa a I_0(\kappa h_s)} \right) \right], \quad (2)$$

206 where F_d^* is a function of the radius a of a sphere (for simplicity we consider bacterial cells to
 207 be spheres); the distance of the center of the sphere to the collector surface, F_d^* , is estimated
 208 to be 1.7; η is the viscosity of the liquid ($\eta = 3.19 \text{ kg m}^{-1} \text{ h}^{-1}$), ε_r is the dielectric constant of
 209 water (78.5), ε_0 ($8.85 \times 10^{-12} \text{ F m}^{-1}$) is the dielectric permittivity in vacuum, ζ_s is the zeta
 210 potential of the sensor surface in the experimental conditions, and E is the electric field strength
 211 applied. I_0 and I_1 are zero-order and first-order modified Bessel functions, and κ^{-1} is the
 212 thickness of the electric double layer. The electrophoretic drag force F_{EP} follows the
 213 Smoluchowski equation (Eq. 3):

$$214 \quad F_{EP} = 6\pi\eta a V_{EP} = 6\pi\eta a \times \frac{2\varepsilon_0 \varepsilon_r \zeta_{bac} E}{3\eta} f(\kappa a), \quad (3)$$

215 where ζ_{bac} is the zeta potential of the bacteria at the given experimental conditions; $f(\kappa a)$ values
 216 approach 1.5 at high electrolyte concentration (i.e., 50 and 100 mM); and $f(\kappa a)$ is close 1.0 at
 217 low ionic strength (i.e., 10 mM) for the bacterial radius $a = 0.6 \text{ }\mu\text{m}$.²² The ratio $|F_{EOF}|/|F_{EP}|$ at
 218 the secondary minimum of the DLVO interaction energy above a collector surface (h_s) is
 219 expressed in Eq. 4:

220
$$\frac{|F_{\text{EOF}}|}{|F_{\text{EP}}|} = \frac{F_d^* \xi}{\frac{2}{3} \xi_{\text{bac}} f(\kappa a)} \left[1 - \frac{2I_1(\kappa h_s)}{\kappa a I_0(\kappa h_s)} \right]. \quad (4)$$

221 Eq. 4 indicates that the $|F_{\text{EOF}}|/|F_{\text{EP}}|$ ratio depends on ζ_{bac} , ζ_s , and the thickness of the electric
 222 double layer κ^{-1} , and is therefore strongly influenced by the ionic strength of the electrolyte.

223

224 QCM-D analyses of bacterial deposition

225 QCM-D is an acoustic method that reflects the amount and viscoelastic properties of an
 226 adhering mass by changes in the resonance frequency (Δf) and energy dissipation (ΔD) of an
 227 oscillating crystal-coated sensor surface.^{36–39 60,61} The change in resonance frequency, Δf , can
 228 be described by the Sauerbrey equation in case of rigid attachment and negative Δf :⁶²

229
$$\Delta f = \frac{-2f_0^2 \Delta m}{A \sqrt{\rho_q \mu_q}} = -C_f \Delta m, \quad (5)$$

230 where f_0 denotes the fundamental resonance frequency, A is the electrode area, ρ_q is the density
 231 of quartz ($\rho_q = 2.648 \text{ g cm}^{-3}$), and μ_q is the shear modulus of quartz ($\mu_q = 2.957 \times 10^{10} \text{ N m}^{-2}$).

232 The $\Delta f/\Delta D$ ratio indicates changes in energy dissipation per coupled unit mass and indicates
 233 the rigidity and attachment strength of bacterial adhesion.^{45,47,63} Typically, bacterial adhesion
 234 leads to a negative frequency shift and a positive dissipation shift. Thus, a low negative $\Delta f/\Delta D$
 235 value indicates the buildup of a dissipative soft and fluid film on the QCM-D sensor. In contrast,
 236 higher negative values of $\Delta f/\Delta D$ indicate a more rigid layer.

237

238 Results

239 Electric field effects and electrolyte effects on the calculated F_{net}

240 In order to approximate the DLVO energy profiles and the electrokinetic forces acting on
 241 bacteria above a sensor surface, the physicochemical properties of the sensor surface and the
 242 bacteria were determined in 10, 50, and 100 mM PB solutions. While the quartz sensor was

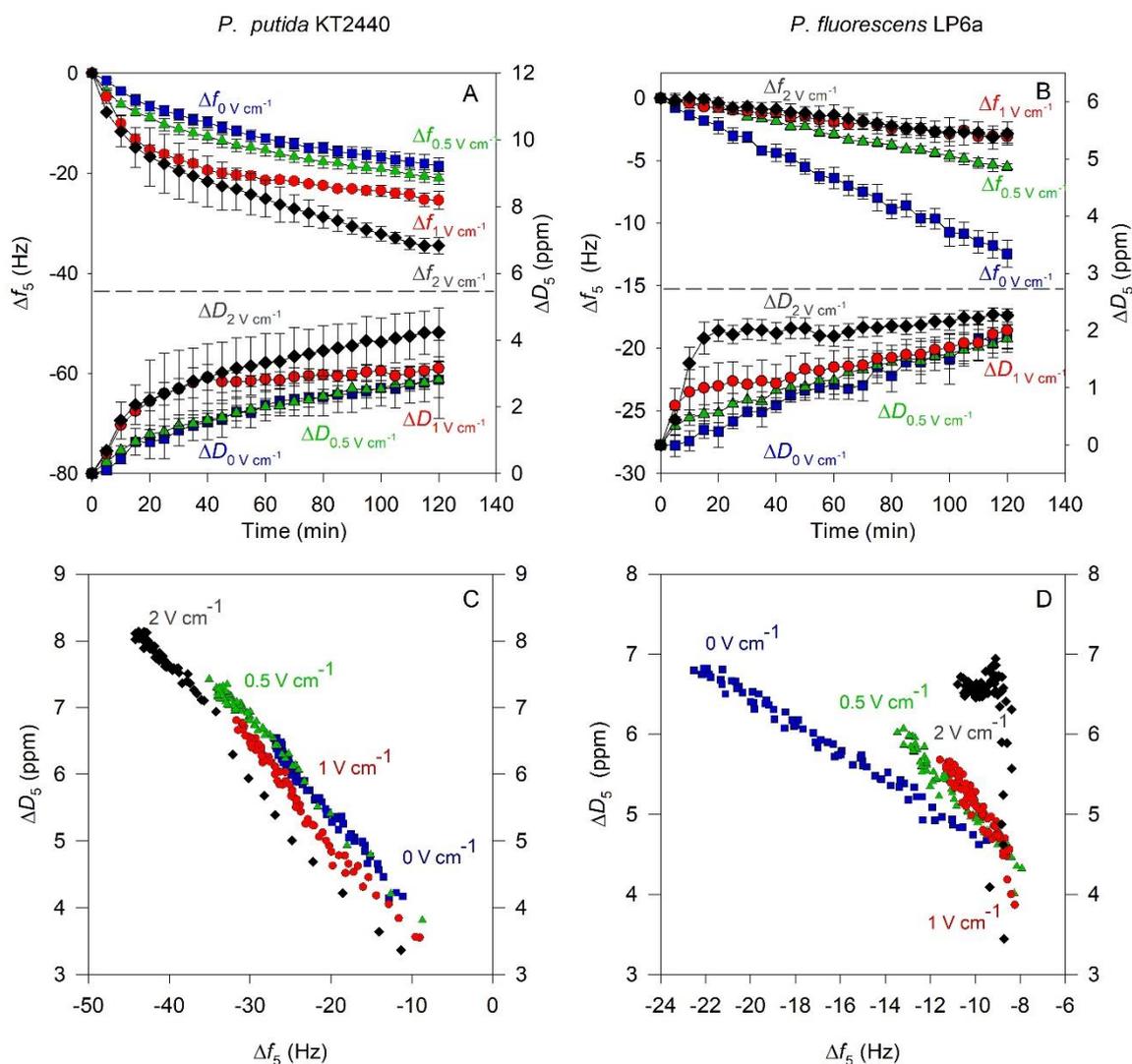
243 hydrophilic (water contact angle, $\theta_w = 21^\circ$), both bacterial strains were moderately
244 hydrophobic ($\theta_{w,KT2440} = 70^\circ$; $\theta_{w,LP6a} = 46^\circ$; Table S1). The sensor surface and both bacterial
245 strains were negatively charged in all PB concentrations (Table 1), with more negative zeta
246 potentials at lower ionic strengths (i.e. shifts from -21 mV (10 mM PB) to -8 mV (100 mM
247 PB) of the sensor, -30 mV to -11 mV (strain KT2440) and -53 mV to -36 mV (strain LP6a)
248 (Table 1). Calculated DLVO interaction energy profiles between the bacteria and the QCM-D
249 quartz sensor surfaces (Fig. S2) all exhibited secondary minima, suggesting reversible bacteria
250 attachment at all PB concentrations. Secondary minima were found at separation distances of
251 $3.2 - 20.6$ nm (Table S2). Corresponding attractive DLVO forces (F_{DLVO}) depended on the ionic
252 strength of the PB and ranged from 0.15 pN (10 mM) to 3.26 pN (100 mM) for strain KT2440
253 and from 0.15 pN (10 mM) to 2.31 pN (100 mM) for strain LP6a (Table 1). Table 1 summarizes
254 the magnitudes of the forces F_{HF} , F_{EOF} , F_{EP} , and F_{net} that we defined as the sum of the
255 magnitudes of F_{HF} , F_{EOF} , and F_{EP} , and F_{DLVO} , disregarding distinct directions of electrokinetic
256 and DLVO forces (Eq. 1). As sensor and bacterial surfaces had negative zeta potentials (Table
257 1), the direction of F_{EP} was opposed to the direction of F_{EOF} , and the magnitudes of F_{EP} were of
258 opposite sign to the magnitudes of F_{EOF} . While the extent of F_{HF} was assumed to be independent
259 of experimental variations, the magnitudes of F_{EOF} and F_{EP} (expressed by $|F_{EOF}|$ and $|F_{EP}|$)
260 increased proportionally to E (Eqs. 2 and 3), and decreased at rising electrolyte concentrations.
261 F_{net} thus depended on the electric field strength and the ionic strength of the PB (Table 1): at
262 any given electric field strength, higher PB concentrations increased the F_{net} of both bacterial
263 strains. At a given ionic strength, however, the F_{net} of the two bacterial strains revealed
264 dissimilar trends at increasing E : in 50 mM and 100 mM PB; an increase in E from 0.5 V cm^{-1}
265 to 2 V cm^{-1} increased F_{net} by ca. 10-20% for strain KT2440 and decreased F_{net} by ca. 700%
266 for strain LP6a (Table 1).

267 **Table 1.** Overview of cell density, zeta potential, and the calculated forces acting on a bacterium (*P.*
 268 *putida* KT2440 or *P. fluorescens* LP6a) at a distance of the secondary minimum in the presence and
 269 absence of a DC electric current at different electrolyte strengths.

270

		<i>P. putida</i> KT2440			<i>P. fluorescens</i> LP6a		
		10 mM	50 mM	100 mM	10 mM	50 mM	100 mM
DLVO force (pN) ^a	F_{DLVO}	0.15	1.45	3.26	0.15	1.43	2.31
Hydraulic shear force (pN) ^b	F_{HF}	0.50	0.50	0.50	0.50	0.50	0.50
Electroosmotic shear force per V cm⁻¹ (pN)	F_{EOF}	3.70	1.90	1.80	3.70	1.90	1.80
Electrophoretic drag force per V cm⁻¹ (pN)	F_{EP}	-3.95	-1.80	-1.45	-6.99	-5.69	-4.74
Net force (pN) ^c							
$E = 0 \text{ V cm}^{-1}$	$F_{Net,ND}$	0.65	1.95	3.76	0.65	1.93	2.81
$E = 0.5 \text{ V cm}^{-1}$	$F_{Net,0.5V \text{ cm}^{-1}}$	0.53	2.00	3.94	-1.00	0.03	1.34
$E = 1 \text{ V cm}^{-1}$	$F_{Net,1V \text{ cm}^{-1}}$	0.40	2.05	4.11	-2.65	-1.86	-0.13
$E = 2 \text{ V cm}^{-1}$	$F_{Net,2V \text{ cm}^{-1}}$	0.15	2.15	4.46	-5.95	-5.65	-3.07
Cell density (10⁶ cells cm⁻²) ^d	d_c						
$E = 0.0 \text{ V cm}^{-1}$	$d_{c, \text{no DC}}$	2.5 ± 0.2	1.8 ± 0.6	4.1 ± 0.6	1.4 ± 0.2	2.9 ± 0.4	3.3 ± 0.3
$E = 0.5 \text{ V cm}^{-1}$	$d_{c, 0.5V \text{ cm}^{-1}}$	1.2 ± 0.3	2.1 ± 0.6	4.8 ± 0.9	1.2 ± 0.3	0.9 ± 0.2	1.9 ± 0.2
$E = 1 \text{ V cm}^{-1}$	$d_{c, 1V \text{ cm}^{-1}}$	1.2 ± 0.2	3.5 ± 0.8	5.0 ± 0.6	0.9 ± 0.1	1.0 ± 0.3	1.3 ± 0.3
$E = 2 \text{ V cm}^{-1}$	$d_{c, 2V \text{ cm}^{-1}}$	0.9 ± 0.4	3.6 ± 0.9	9.9 ± 2.1	0.7 ± 0.1	0.8 ± 0.2	1.1 ± 0.3
Zeta potential (-mV)							
Bacteria	ζ_{bac}	-30 ± 1	-14 ± 2	-11 ± 1	-53 ± 2	-43 ± 2	-36 ± 3
		Sensor surface					
Silica ^e	ζ_s	-21 ± 2	-12 ± 1	-8 ± 1			

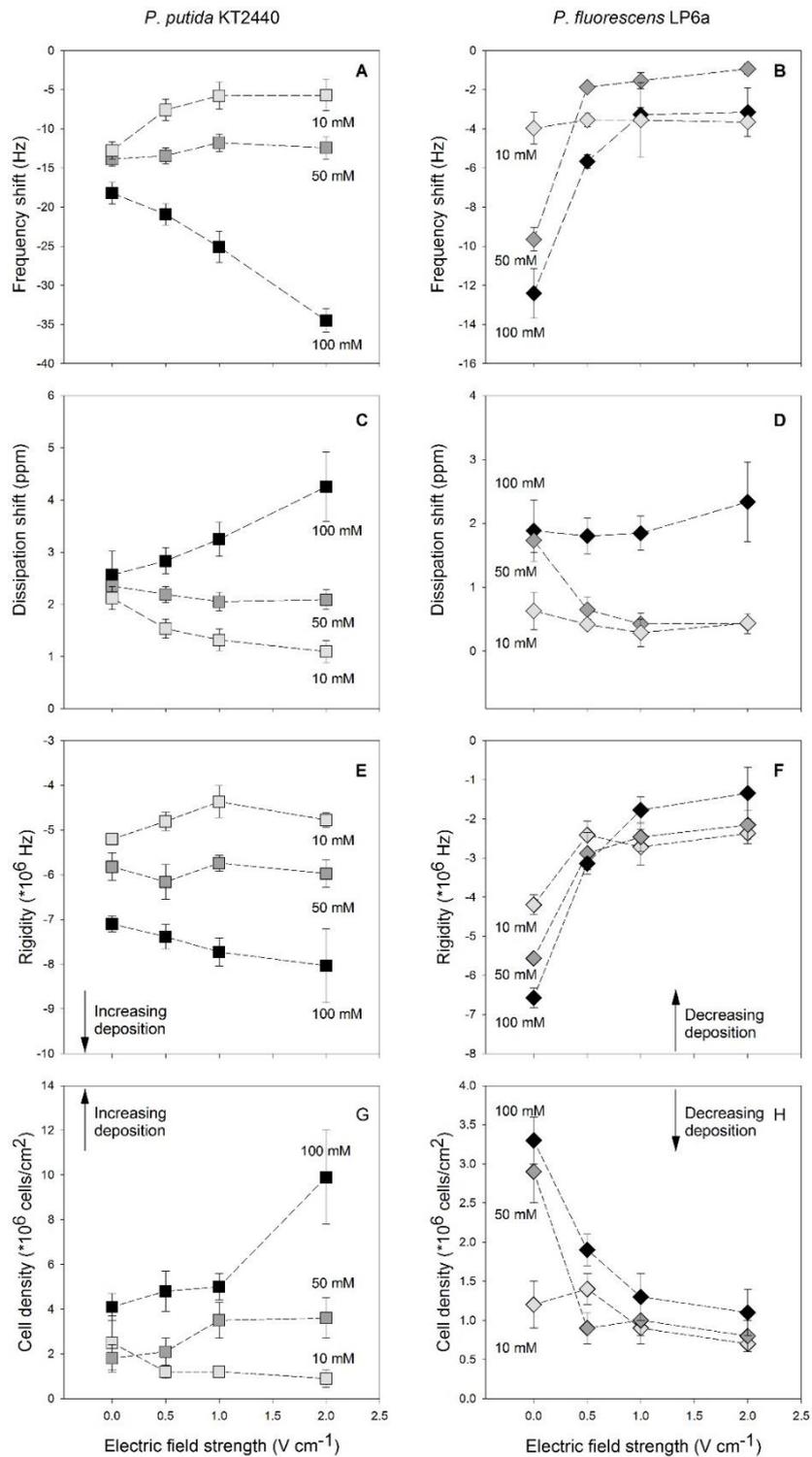
271 ^a For calculation cf. Eq. S10; ^b F_{HF} calculated for flow velocity of $6 \times 10^{-7} \text{ m s}^{-1}$ (cf. Eq. S11); ^c cf. Eq. 1; ^d Microscopically
 272 determined cell density after 2 h; ^e Silica sensor surface.



273

274 **Figure 1.** Time dependent frequency shifts (Δf_5) and dissipation shifts (ΔD_5) of *P. putida* KT2440 (Fig.
 275 1A) and *P. fluorescens* LP6a (Fig. 1B) at overtone 5 in 100 mM PB and electric field strengths of $E =$
 276 0 V cm^{-1} (blue squares), $E = 0.5 \text{ V cm}^{-1}$ (green triangles), $E = 1.0 \text{ V cm}^{-1}$ (red circles), and $E = 2.0 \text{ V cm}^{-1}$
 277 1 (black diamonds). Error bars denote the standard deviation of the mean ($n = 3$). Data above and below
 278 the dashed line refer to Δf_5 (left y-axis) and to ΔD_5 (right y-axis), respectively. Panels C and D correlate
 279 the time dependent ΔD_5 and Δf_5 of *P. putida* KT2440 and *P. fluorescens* LP6a.

280



281

282

283 **Figure 2.** Effect of the electric field strength on the frequency shift (Δf_5 ; Figs. 2A and 2B), the
 284 dissipation shift (ΔD_5 ; Figs 2C and 2D), the rigidity of bacterial attachment ($\Delta f_5/\Delta D_5$, Figs. 2E and 2F),
 285 and the cell density on the sensor surface (Figs. 2G and 2H). Bacterial deposition of *P. putida* KT2440,
 286 Figs. 2A, 2C, 2E, and 2G, and *P. fluorescens* LP6a, Figs. 2B, 2D, 2F, and 2H, after two hours (cf. Fig.
 287 S3) at overtone 5 in 10 mM (light gray), 50 mM (dark gray) and 100 mM (black) PB.

288 **Electric field and electrolyte effects on Δf and ΔD and derived cell attachment rigidity**

289 QCM-D experiments recorded frequency shifts and dissipation shifts at overtones 1, 3, 5, 7, 9,
290 11, and 13 (Fig. S3) during 120 minutes of bacterial deposition. While overtone 1 was poorly
291 stable and overly sensitive, all other overtones showed similar trends (Figs. S4 and S5). In the
292 following, we analyze and discuss overtone 5 as a representative signal using the
293 frequency/dissipation baseline in cell-free PB as a reference to calculate the frequency and
294 dissipation shifts of the bacteria deposition (Figs. 1, S4, and S5). Figure 1 exemplifies Δf_5 and
295 ΔD_5 shifts of both strains in 100 mM PB at electric field strengths of $E = 0, 0.5, 1, \text{ and } 2 \text{ V cm}^{-1}$
296 ¹. Here, pumping bacteria over the sensor surface resulted in a decrease in frequency shifts and
297 an increase in dissipation shifts; Δf_5 and ΔD_5 varied at different experimental conditions (Figs.
298 1A, 1B, S4, and S5). Generally, the rates of Δf_5 and ΔD_5 were higher at the beginning (0-15
299 minutes) of bacterial deposition than at the end of bacterial deposition (Figs. 1A, 1B for 100
300 mM PB and Figs. S4, S5 for 10 and 50 mM PB), while the ratio $\Delta f_5/\Delta D_5$, an indicator of
301 attachment rigidity, generally exhibited a linear correlation with Δf_5 and ΔD_5 ranges, with
302 coefficients of determination (r^2) of > 0.95 (Figs. 1C, 1D, Table S4). Figs. 2A-F summarize
303 Δf_5 , ΔD_5 , and $\Delta f_5/\Delta D_5$ ratios at the end of the deposition experiments. While signals of strains
304 KT2440 and LP6a differed depending on the experimental conditions, the effects observed
305 were proportional to the electric field strength applied; i.e., a higher voltage resulted in stronger
306 observed effects. For strain KT2440 in 100 mM PB, for instance, Δf_5 decreased from -18.2 Hz
307 ($E = 0 \text{ V cm}^{-1}$) to -34.5 Hz ($E = 2.0 \text{ V cm}^{-1}$) while ΔD_5 increased from 2.56 ppm to 4.25 ppm
308 (Figs. 2A and 2 C). Such shifts resulted in clear increases in the calculated rigidity (i.e., more
309 negative $\Delta f_5/\Delta D_5$ ratios; Fig. 2E). In contrast, Δf_5 , ΔD_5 , and $\Delta f_5/\Delta D_5$ ratios of strain LP6a in
310 100 mM PB increased with increasing electric field strengths; i.e., Δf_5 from -12.4 Hz to -3.14
311 Hz, ΔD_5 from 1.89 ppm to 2.34 ppm, and $\Delta f_5/\Delta D_5$ from -6.56 to -1.34 MHz (Figs. 2 B, D, F).
312 Decreasing PB concentrations from 100 mM to 10 mM resulted in lower shifts of Δf_5 , ΔD_5 ,

313 and $\Delta f_5/\Delta D_5$ in DC free controls and smaller DC-induced changes, respectively. For strain LP6a,
314 an electric field as weak as $E = 0.5 \text{ V cm}^{-1}$ resulted in distinct changes in Δf_5 , ΔD_5 , and $\Delta f_5/\Delta D_5$
315 at all PB concentrations. In contrast, DC field effects on the trends of Δf_5 , ΔD_5 , and $\Delta f_5/\Delta D_5$ of
316 strain KT2440 varied with the concentration of the electrolyte. At PB concentrations of 10 and
317 50 mM, DC fields decreased the rigidity of attached KT2440 cells, while more negative
318 $\Delta f_5/\Delta D_5$ ratios (i.e., more rigid attachments) were observed at increasing E .

319 **Electric field and electrolyte effects on cell density of attached bacteria**

320 The number of cells attached to the sensor surface was counted microscopically at the end of
321 the deposition experiments. The cell density (d_c) and the surface coverage (cf. Eq. S12) of cells
322 attached to the quartz sensor surface (1.54 cm^2) were approximated. The d_c varied from $0.9 \times$
323 10^6 to $9.9 \times 10^6 \text{ cells cm}^{-2}$ (strain KT2440) and from 0.7×10^6 to $3.3 \times 10^6 \text{ cells cm}^{-2}$ (strain
324 LP6a) (Table 1; Figs. 2G and 2H). This corresponds to maximal coverages of the sensor surface
325 (Table S3) of 5.5% and 1.8%, respectively. In strain LP6a (where d_c at 10 mM and 50 mM
326 were similar), the cell density increased in the order of $d_c (10 \text{ mM}) < d_c (50 \text{ mM}) < d_c (100$
327 $\text{mM})$ at all electric field strengths (Table 1). At a given PB concentration, however, the strength
328 of the electric fields evoked distinct d_c differences between the two bacterial strains (Table 1
329 and Figs. 2E and 2F). An increase in E resulted in a decrease in the d_c of strain LP6a at all
330 electrolyte concentrations, suggesting that DC electric fields reduced the deposition of LP6a
331 cells to the sensor surface even at weak E . For strain KT2440 however, an increase in the
332 electric field decreased cell attachment to the sensor in 10 mM PB, but promoted cell
333 attachment in 50 mM and 100 mM PB (Table 1 and Figs. 2E and 2F). Cell density data for both
334 bacterial strains thereby showed similar relative trends in Δf_5 and ΔD_5 (Figs. 2A, 2B).

335

336

337 **Discussion**

338 **Assessment of DC-induced deposition effects by QCM-D monitoring**

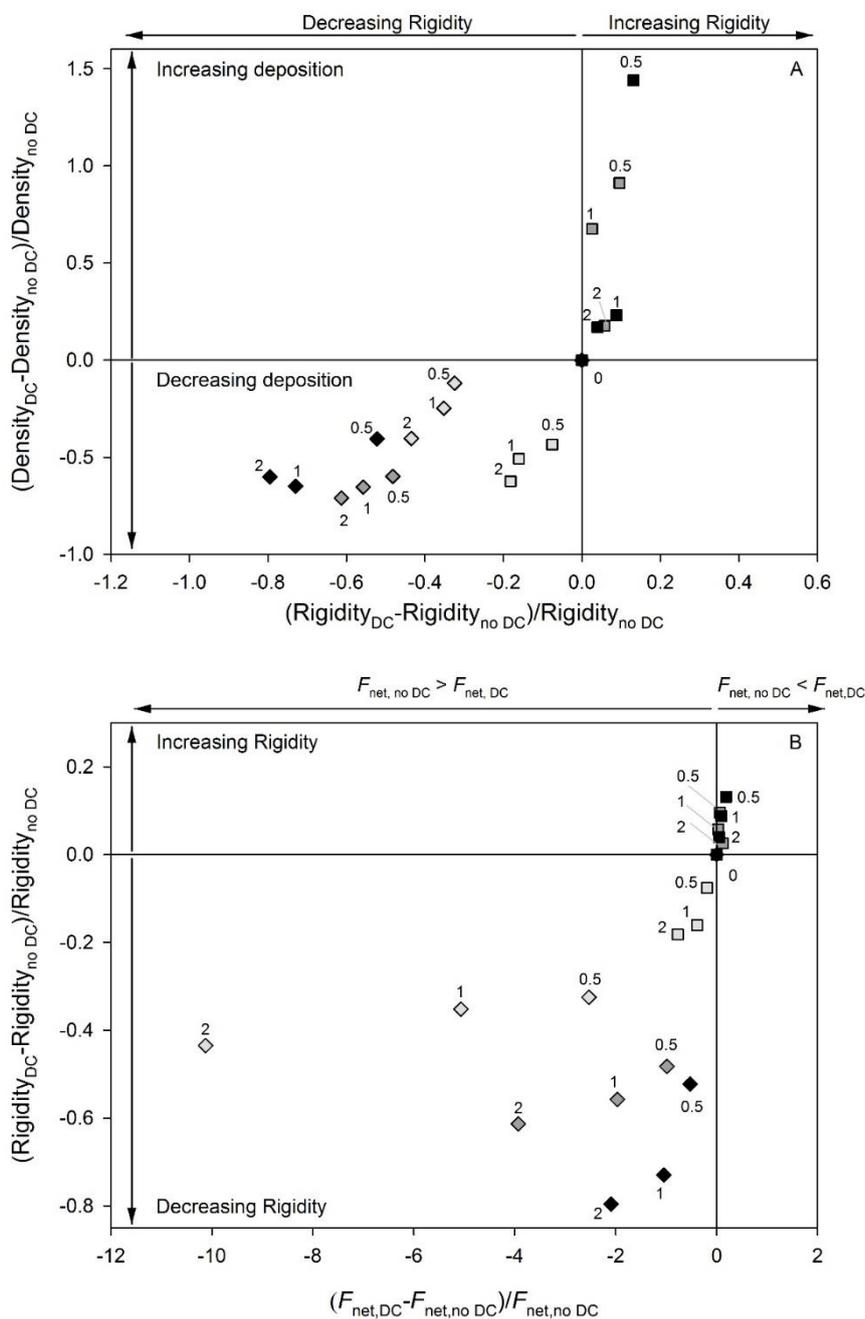
339 Motivated by recent work that suggested that bacterial deposition and transport in percolation
340 systems is influenced by electrokinetic forces,¹⁴ we studied DC electric field effects on
341 bacterial deposition using real-time QCM-D monitoring at varying PB concentrations (10-100
342 mM) and electric field strengths (0-2 V cm⁻¹). Electrolyte concentration and electric field
343 strength are key drivers of electrokinetic shear and drag forces acting on bacteria. QCM-D
344 signals were further compared to cell density. The results are discussed based on
345 approximations of the net force (F_{net} ; Eq. 1) acting on a bacterium at the distance of reversible
346 attachment (i.e., at the secondary minimum of the DLVO interaction energy of bacterial
347 adhesion, G_{DLVO} , Eq. S1 and Fig. S2). Except for strain LP6a at 2 V cm⁻¹, we found good
348 correlation between the resonance frequency (Δf_5) and the dissipation energy (ΔD_5) in bacterial
349 strains KT2440 and LP6a in all experiments (Figs. 1C and 1D). Based on work by Gutman et
350 al.,⁴⁵ we used $\Delta f_5/\Delta D_5$ ratios to indicate attachment rigidity⁴² and cell deposition.³⁴ Our data
351 showed good correlation between $\Delta f_5/\Delta D_5$ and the microscopically determined cell density (d_c).
352 (Fig. 3A). Backed by both the attachment rigidity and the cell density, we found that weak DC
353 fields clearly changed the deposition patterns of strains KT2440 and LP6a compared to DC-
354 free controls (Figs. 2E-2H). Observed deposition effects were proportional to the electric field
355 strength applied (i.e., stronger effects were exhibited at higher E), yet were dependent on the
356 bacterial cell surface properties and the PB ionic strength (Fig. 2).

357 **Prediction of DC-induced bacterial deposition effects**

358 According to the Derjaguin, Landau, Verwey, and Overbeek (DLVO) theory,⁵³ deposition of a
359 bacterium to a sensor surface requires that the net kinetic energy of the bacterium is lower than
360 the DLVO interaction energy at the distance of reversible attachment.^{56,64} Prediction of DC

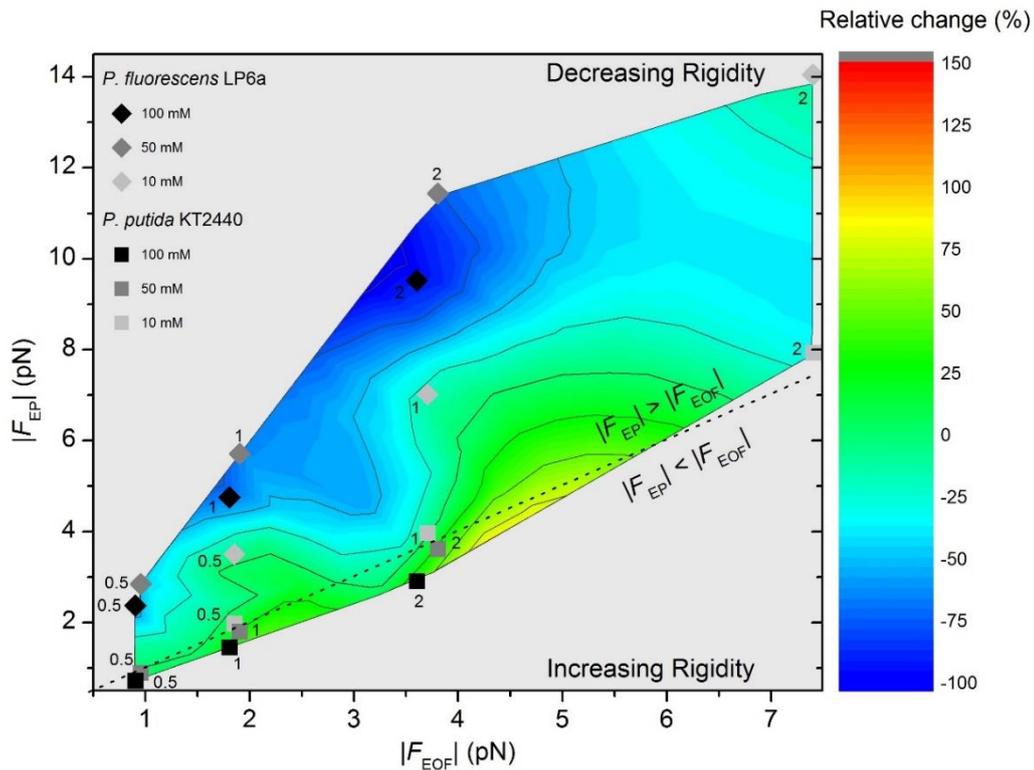
361 electric field effects on bacterial deposition should therefore consider additional electrokinetic
362 forces acting on depositing cells; for example, electroosmotic shear and electrophoretic drag
363 forces have powerful effects on the movement of bacteria and (bio-)colloidal particles.^{7,9,16} We
364 correlated DC-induced deposition effects with F_{net} shifts (Figs. 3B and S7); i.e., the attachment
365 rigidity ($\Delta f_5/\Delta D_5$) and the cell density (d_c) were correlated with the F_{net} acting on a bacterium
366 at the secondary minimum above the sensor surface. For easier comparison, all data were
367 normalized for DC-free controls, using $((\Delta f_5/\Delta D_5)_{\text{DC}} - (\Delta f_5/\Delta D_5)_{\text{no DC}})/(\Delta f_5/\Delta D_5)_{\text{no DC}}$ i.e.: for
368 attachment rigidity, $(d_{c,\text{DC}} - d_{c,\text{no DC}})/d_{c,\text{no DC}}$ for cell density, and $(F_{\text{net,DC}} - F_{\text{net,no DC}})/F_{\text{net,no DC}}$
369 for normalized net force shifts, respectively. In doing so, we found good correlation between
370 the normalized d_c and QCM-D derived rigidity (Fig 3A) at all electric field strengths and buffer
371 concentrations tested. Increasing attachment rigidity was mirrored by higher d_c , while
372 decreasing attachment rigidity resulted in lower d_c (Fig. 3A). This highlights QCM-D as a
373 useful approach to assess and predict the influence of DC electric fields on bacterial deposition:
374 At $F_{\text{net,DC}} > F_{\text{net,noDC}}$, increased attachment rigidity (Fig 3B) and higher d_c (Fig. S7) were
375 observed, and at $F_{\text{net,DC}} < F_{\text{net,noDC}}$, lower attachment rigidity (Fig 3B) and lower d_c (Fig. S7)
376 were observed. As F_{EOF} and F_{EP} are of opposite sign in our experimental system, their relative
377 strengths are a driver of $F_{\text{net,DC}}$ (Eq. 1) and, thus, of observed electrokinetic effects on bacterial
378 deposition (Figs. 4 and S8). If $|F_{\text{EOF}}| > |F_{\text{EP}}|$, DC fields promote attachment rigidity and d_c and
379 *vice versa*, respectively.¹⁴ Therefore, $|F_{\text{EOF}}|/|F_{\text{EP}}|$ was a good predictor of bacterial electrokinetic
380 effects on cell attachment rigidity and bacterial deposition in all conditions tested. The heat
381 maps in Figs. 4 and S8 show the effects of $|F_{\text{EOF}}|$ and $|F_{\text{EP}}|$ on the normalized DC-induced
382 rigidity and d_c changes. They reveal the importance of $|F_{\text{EP}}|$ for cell deposition at a given $|F_{\text{EOF}}|$,
383 independent of bacterial strain, electrolyte strength, and applied electric field. The high degree

384 of convergence between changes in rigidity and changes in d_c further indicates that QCM-D is
 385 a useful and expedient tool for the real-time analysis of electrokinetic deposition.



386
 387 **Figure 3.** Correlation of normalized changes in DC-induced cell density, rigidity of cell attachment
 388 (Fig. 3A), DC-induced net force ($F_{\text{net,DC}}$, cf. Eq. 1), and rigidity of cell attachment (Fig. 3B), respectively.
 389 All plots reflect data after two hours of deposition of *P. putida* KT2440 (squares) and *P. fluorescens*
 390 LP6a (diamonds) exposed to PB at concentrations of 10 mM (light gray), 50 mM (dark gray), and 100
 391 mM (black), and DC electric field strengths of $E = 0, 0.5, 1.0, \text{ or } 2.0 \text{ V cm}^{-1}$ (cf. digits at the symbols).

392



393

394 **Figure 4.** Calculated effects of the electroosmotic shear force $|F_{\text{EOF}}|$ and the electrophoretic drag force
 395 $|F_{\text{EP}}|$ on DC-induced normalized changes in the rigidity of cell attachment after two hours of deposition
 396 of *P. putida* KT2440 (squares) and *P. fluorescens* LP6a (diamonds). Experiments were performed in
 397 PB at concentrations of 10 mM (light gray), 50 mM (dark gray), and 100 mM (black), and DC electric
 398 field strengths of $E = 0, 0.5, 1.0,$ or 2.0 V cm^{-1} (cf. digits at the symbols). Data points above the dashed
 399 line (i.e. $|F_{\text{EP}}| > |F_{\text{EOF}}|$) and below the dashed line (i.e. $|F_{\text{EP}}| < |F_{\text{EOF}}|$) refer to decreased and increased
 400 rigidity, respectively, compared to DC-free controls.

401

402 **Relevance of findings to environmental applications**

403 Electrokinetic transport processes are often applied in civil and environmental engineering such
 404 as for wood preservation or for contaminant removal. As an alternative to physical filtration,
 405 electrokinetic approaches can be used to pre-concentrate large molecules and nanoparticles
 406 using the double layer properties of nanochannels (“electrokinetic trapping”⁶⁵). Here we applied
 407 electrokinetic forces to influence bacterial deposition on surfaces. Electrokinetic deposition
 408 approaches may be used in future applications to retain unwanted bacteria in drinking water
 409 purification systems or - vice versa - to reduce bacterial deposition and subsequent bio-fouling
 410 in engineered systems. . The relative strength of F_{EOF} and F_{EP} acting on bacteria at a distance of

411 the secondary DLVO minimum above a surface was found to be a good predictor of
412 electrokinetic effects on cell deposition. According to Eq. 4, the $|F_{\text{EOF}}|/|F_{\text{EP}}|$ ratio is influenced
413 by the electric field strength, the ionic strength of the electrolyte, the zeta potentials of bacteria
414 and bacteria collector surfaces, and the thickness of the electric double layer. QCM-D allows
415 for fast, real-time and accurate high throughput monitoring of bacterial deposition by easily
416 changing the drivers of the $|F_{\text{EOF}}|/|F_{\text{EP}}|$ ratio. It can be used to predict electrokinetic effects on
417 bacterial deposition in environmental and biotechnological applications (e.g., elimination of
418 unwanted bacteria in drinking water or the prevention of biofilm induced corrosion).
419 Knowledge on DC-effects also allows to manage electrokinetic bacterial dispersal in
420 subsurface porous media and e.g. to change microbial community structures and functions and
421 to promote contaminant biodegradation in disturbed ecosystems.^{66,67} Electrokinetic effects may
422 also improve the transport of nutrients by electromigration or change the interactions of
423 contaminants with sorbents,^{68,69} thereby enhancing the biodegradation of contaminants during
424 engineered clean-up of contaminated soil or waters.

425 **Supporting Information.** The SI contains 4 tables and 8 figures as well as calculations of the
426 DLVO interaction force between bacteria and a solid surface (F_{DLVO}), and the hydraulic shear
427 force F_{HF} , resp. It further describes the estimation of the bacterial coverage of attached bacterial
428 cells on the sensor and provides a code for ImageJ automatic cell counting of images taken
429 with a Hemacytometer.

430 **Acknowledgments.** This work was performed in the frame of the Helmholtz Alberta Initiative
431 and contributes to the research program topic CITE of the Helmholtz Association. We
432 acknowledge financial support by the China Scholarship Council (CSC) and the German
433 Academic Exchange Service (DAAD). The authors thank Dr. Luis Rosa for helpful discussions
434 and Jana Reichenbach, Rita Remer, and Birgit Würz for skilled technical help.

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